

Locking Ideas

ede, 17.11.02

1 Notations

ω_{ext} = frequency of external signal, ω_{free} = frequency of free running self-pulsation, $\delta = \omega_{free} - \omega_{ext}$ detuning between both.

φ = phase shift between the self-pulsation (pulses at the internal wavelength) and the externally injected pulses at an arbitrary but fixed position of the device, e.g., at one facet.

Hence, $\dot{\varphi}$ = instantaneous frequency of the internal self-pulsation (relative to ω_{ext}).

2 Assumptions and Conclusions

Assumption 1: The instantaneous frequency $\dot{\varphi}$ is a unique function of φ : $\dot{\varphi} = \omega(\varphi)$.

This makes sense at least for sufficiently small detuning and signal intensity. It is obviously true in the limit of vanishing signal intensity, when holds $\omega(\varphi) = \delta = \text{const.}$ With increasing signal intensity, $\omega(\varphi)$ will show more and more variation.

Conclusion 1: Obviously, $\omega(\varphi)$ is a periodic function.

Conclusion 2: Locking appears if and only if $\omega(\varphi) = 0$ has a solution.

Conclusion 3: The locking dynamics is governed by the equation

$$\dot{\varphi} = \omega(\varphi). \quad \text{General solution:} \quad t = \int_{\varphi_1}^{\varphi_2} \frac{d\varphi}{\omega(\varphi)}, \quad (1)$$

where t is the time needed for the phase change from the initial value φ_1 to the final value φ_2 .

Qualitative discussion: If $\min(\omega(\varphi)) > 0$, there is no locking. But the phase φ will change slowliest around the minimum of $\omega(\varphi)$. If this minimum is closely above 0, φ will become nearly stationary for a comparatively long time. Such φ -plateaus did we observe in the numerics, they are a precursor of locking. The slope of the plateaus is given by the minimum of $\omega(\varphi)$. The border of the locking range is achieved when $\min(\omega(\varphi)) = 0$. Beyond, we have at least one zero of $\omega(\varphi)$, say at φ_0 . It is approached exponentially fast with a time constant $1/\omega'(\varphi_0)$, which is very long at the border of the locking range but smaller in its interior. To be concrete, I consider now a model.

$$\textbf{Model:} \quad \omega(\varphi) = \delta + p \sin(\varphi). \quad \implies \quad \boxed{\dot{\varphi} = \delta + p \sin(\varphi)} \quad (2)$$

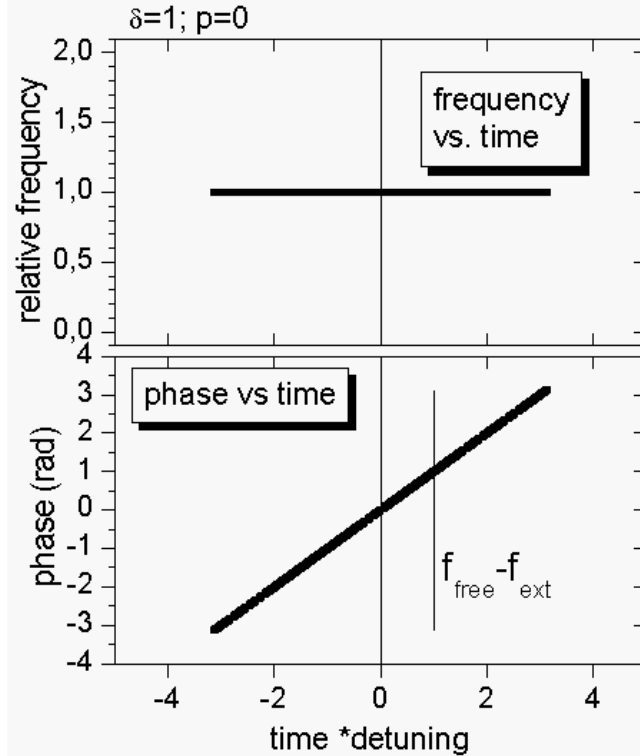
This seems to be Adler's equation, if I remember right. It can be solved analytically:

$$t = \frac{2}{\sqrt{\delta^2 - p^2}} \arctan \frac{\delta \tan(\varphi/2) + p}{\sqrt{\delta^2 - p^2}} \quad (p < \delta) \quad (3)$$

$$t = \frac{1}{\delta} \tan \left(\frac{\varphi}{2} - \frac{\pi}{4} \right) \quad (p = \delta) \quad (4)$$

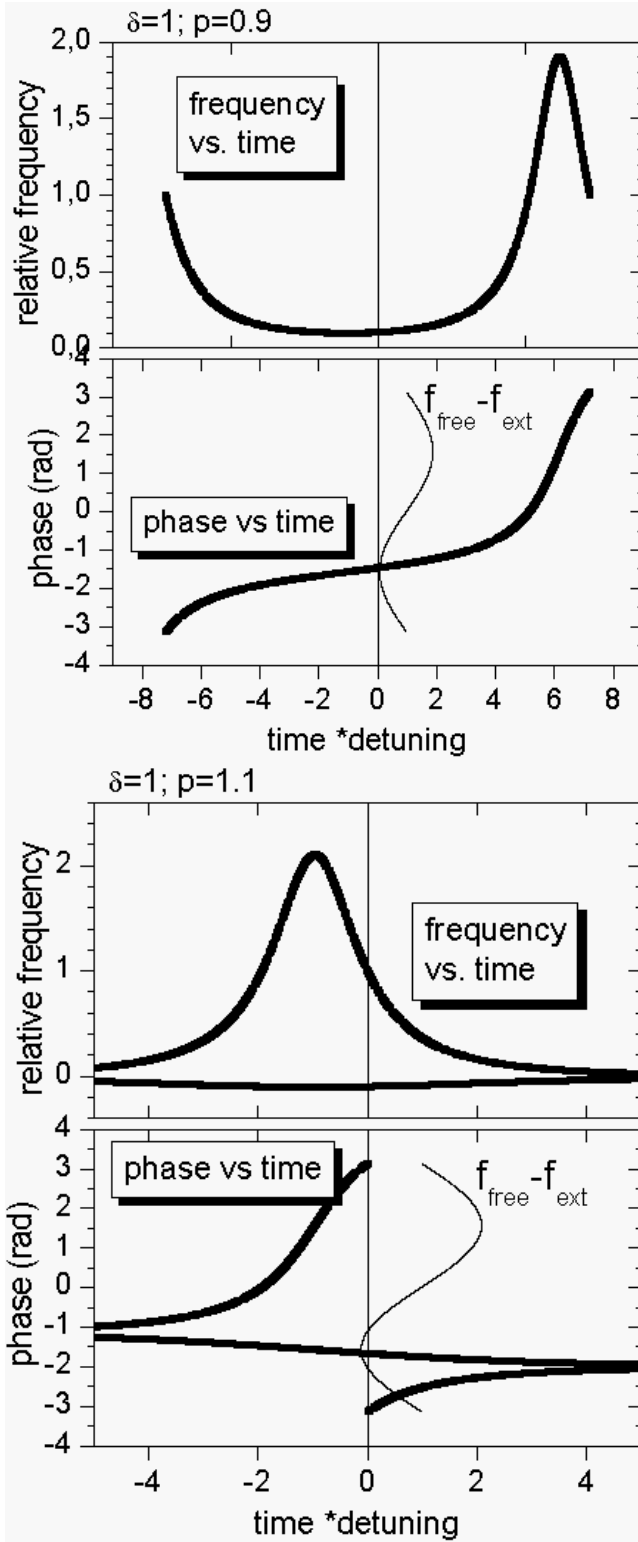
$$t = \frac{1}{\delta} \ln \left| \frac{\delta \tan(\varphi/2) + p - \sqrt{p^2 - \delta^2}}{\delta \tan(\varphi/2) + p + \sqrt{p^2 - \delta^2}} \right| \quad (5)$$

I've plotted these solutions for some characteristic cases, assuming $\delta = 1$. Abscissa is always the dimensionless time δt . Lower panels: thick: φ versus δt , thin: φ vs. $\omega(\varphi)$, i.e., the model (2), to indicate the position of the zeros φ_0 .



Case p=0.

Without signal, the frequency remains constant, the phase raises linearly. The time for one phase period (change by 2π) is just $2\pi/\delta$.

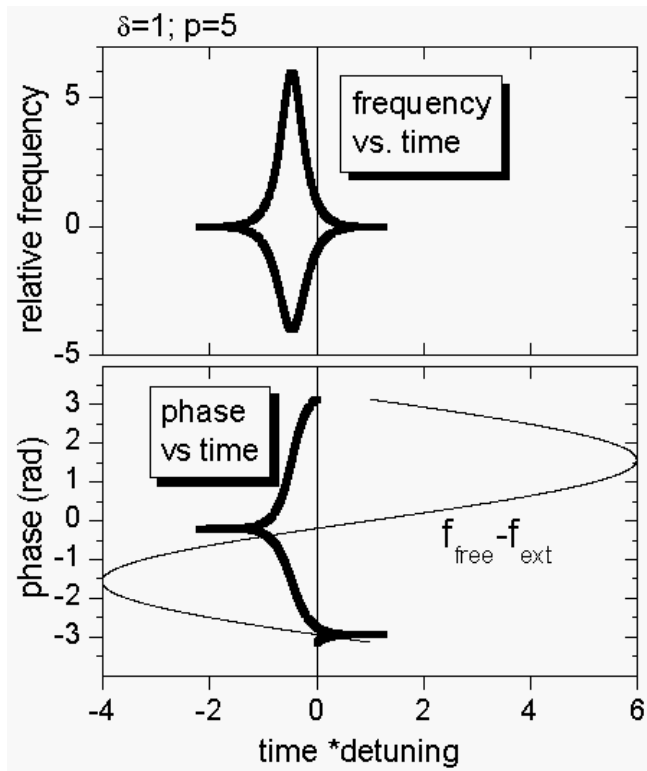


Case $p=0.9$.

There is no zero yet of (2), we are still outside the locking range. However, the minimum $\min(\omega(\varphi)) = \delta - p = 0.1$ is already close to zero (cf. thin line in the lower panel). Hence, we clearly observe a range with stagnating phase, the precursor of locking known from the numerics with LDSL. Note that the time period is enhanced now to approximately $14/\delta$ due to this stagnation.

Case $p=1.1$

Now we have two zeros of $\omega(\varphi)$. This one with a negative slope is the stable one, which the phase approaches. Depending on the starting phase, the time to lock is comparable with the extension of the time axis, roughly $10/\delta$. We are still close to the border of the locking range.



Case $p=5$

Now we are more in the middle of the locking range, the time needed is only roughly $1/\delta$ yet.