

Correlation function and time from power spectra

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After having read Born-Wolff chapter 10.2, I came to conclusions very similar to Oleg's brief introduction. Here a summary.

I. PRELIMINARIES

Be $P(t)$ a transient of optical intensities detected by the photodiode at the ESA.

Be $S(\nu)$ the power spectrum (spectral density) of $P(t)$, defined as [1]

$$S(\nu) \stackrel{\text{def}}{=} \lim_{T \rightarrow \infty} \frac{|\overline{F_T(\nu)}|^2}{2T} \stackrel{(WK)}{=} \int_{-\infty}^{\infty} d\tau G(\tau) e^{i\omega\tau}, \quad G(\tau) = \langle P(t)P(t+\tau) \rangle, \quad \omega = 2\pi\nu. \quad (1)$$

$F_T(\nu)$: Fourier transform of $P(t)$ truncated outside the interval $-T < t \leq T$, (WK): Wiener-Khinchin-Theorem [formula (28) in Born-Wolff], $|\overline{\dots}|^2$: ensemble average [2], $\langle f(t) \rangle$: temporal average over t [3].

Be $P(\nu_m)$ the spectrum measured with the ESA on an equidistant grid ν_m of frequencies.

Assumption 1: $P(\nu_m)$ is a reasonable approximation of the powerspectrum of $P(t)$.

Nothing else makes sense.

Conclusion: $P(\nu_m)$ must be quadratic in the power transient $P(t)$.

This conclusion is easily checked: look to a spectrum with at least one peak well above noise background. Attenuate the input signal by x dB. Then measured spectrum should go down by 2x dB. I'm sure that one of the usable scales (hopefully dBm) fulfills this requirement [4].

Assumption 2: $P(\nu_m)$ represents the contribution of $S(\nu)$ to an grid interval $\nu_m - \nu_{m-1}$ around ν_m , normalised to a certain resolution bandwidth $\Delta\nu_{\text{res}}$.

The corresponding formula is

$$P(\nu_m) = \frac{\Delta\nu_{\text{res}}}{\Delta\nu_{\text{grid}}} \int_{-\Delta\nu_{\text{grid}}/2}^{\Delta\nu_{\text{grid}}/2} S(\nu_m + \nu') d\nu'. \quad \Delta\nu_{\text{grid}} = \nu_m - \nu_{m-1}. \quad (2)$$

It seems to me the only reasonable choice by the following reasons:

- Sharp lines in between the frequency grid points should not be lost. Thus, I must integrate over all intervals between.
- The magnitude of the spectrum should not depend on the choosen grid width. Thus, I must average over the grid intervals.
- The used unit is "mW", thus I must multiply the average with some standard frequency interval. I guess, it is the resolution band width, which can be choosen independently of the grid step.

II. CORRELATION FUNCTION

The (nonnormalised) correlation function is

$$C(\tau) \stackrel{def}{=} G(\tau) - \langle P(t) \rangle^2 = \int_{-\infty}^{\infty} S(\nu) e^{-i\omega\tau} d\nu - \langle P(t) \rangle^2. \quad (3)$$

Thus, it can be calculated from the power spectrum. Since the power spectrum is symmetric in ν , this is a real quantity. Only how to calculate the subtracted second term from the power spectra, I don't know at present. However, except in ideal coherent cases it is just $G(\tau \rightarrow \infty)$, i.e. it can be taken from the large time limit of G .

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- [1] I follow chapter 10.2 of Born-Wolff.
 - [2] this corresponds to averaging over many repetitions of the measurement, which can be done with the ESA. Thus, $S(\nu)$ corresponds to the limit of infinitely many repetitions.
 - [3] i.e. $\langle f(t) \rangle = \lim_{T \rightarrow \infty} \int_{-T}^T f(t + t') dt' / 2T$
 - [4] I guess, the optical power P is linearly transferred to a photo current I , which in turn is linearly related to the voltage U at some internal resistivity (50Ω ?) analysed by the electronics of the ESA. The output with the choicable unit "dBm" is probably related to the electrical power consumed by the resistivity, i.e. $\sim U^2 \sim I^2 \sim P^2$, quadratic with the detected optical intensity. Let's look to the manual.